# Ab Initio Calculations of the Geometries and IR Spectra of Two Derivatives of Tetramethyleneethane 

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Following the observation of a triplet EPR signal for tetramethyleneethane (TME) I in 1970, ${ }^{1}$ a conflict between experiment and theory regarding the multiplicity of the ground state of I has arisen. The ab initio calculations of Borden and Davidson ${ }^{2}$ predict a singlet ground state, but Dowd, Chang, and Paik ${ }^{3}$ have found that I obeys the Curie law, supporting the assignment of a triplet

I

II

III
ground state to I. The divergence between theory and experiment might have arisen because I has a perpendicular or twisted conformation, while the early theoretical work was focused on the planar diradical. ${ }^{2}$ However, Du and Borden ${ }^{4}$ have recently predicted, on the basis of ab initio calculations, that I should have a singlet ground state regardless of the angle of twist about the central carbon-carbon bond; they also concluded that the singlet state adopts the perpendicular ( $D_{2 d}$ ) structure. At their optimized structure the singlet-triplet separation is -0.1 to $2.8 \mathrm{kcal} / \mathrm{mol}$, depending on the level of theory applied.

The conflict between theory and experiment has made it desirable to prepare diradicals with unequivocal geometries. Dowd, Chang, and Paik ${ }^{5}$ have recently demonstrated that the 2,3 -di-methylenecyclohexa-1,3-diene diradical II $^{6}$ also exhibits linear Curie law behavior. It was suggested on the basis of an MM2 calculation that this ground-state triplet molecule, with four $\mathrm{sp}^{2}$-hybridized carbon atoms in a six-membered ring, should be nearly planar. ${ }^{5}$ However, the MM2 procedure as currently parameterized is not ideally suited for examining the question of nonplanarity of conjugated diradical systems. This motivated us to calculate the geometry of the triplet state of II as well as that of III by using the ab initio Hartree-Fock procedure. ${ }^{7}$ Roth and co-workers ${ }^{8}$ have concluded, on the basis of its Curie law behavior, that the planar 2,2-dimethyl-4,5-dimethylene-1,3-cyclopentadiene diradical has a triplet ground state.

The structures of both II and III were fully optimized by using the unrestricted Hartree-Fock (UHF) procedure and the 3-21G basis set. ${ }^{9}$ The resulting structures are shown in Figure 1. The calculations yield a $C_{2}$ structure for II, with a dihedral angle of $25.0^{\circ}$ between the planes of the allyl groups. In agreement with UHF calculations with the less flexible STO-3G basis set, ${ }^{10}$ we found that the lowest triplet state of III is planar with $C_{2 v}$ symmetry. The structures of both II and III were shown to be true

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Figure 1. The optimized 3-21G geometries of the diradicals 2,3 -di-methylenecyclohexa-1,3-diene (II) and 2,3-dimethylenecyclopenta-1,3diene (III).
energy minima by calculating the force constants and normal mode frequencies with analytical second derivatives.

The normal mode frequencies, reduced by $10 \%$ to correct roughly for the deficiencies of the HF procedure and for the effects of anharmonicity, are summarized in Table I. The calculated frequencies and IR intensities of II are compared with the experimental spectrum (Figure 2) obtained by Roth et al. ${ }^{69}$ following irradiation at 254 nm of 5,6-dimethylene-2,3-diazabicyclo-[2.2.2]oct-2-ene in an argon matrix. The bands which grew in upon irradiation were identified with those of the diradical II. Our theoretical results are in fairly good agreement with the results reported in the $600-1000-\mathrm{cm}^{-1}$ region of the spectrum. ${ }^{6}$ Likewise, we predict a strong band at $3011 \mathrm{~cm}^{-1}$ in agreement with experiment. We also predict eight fairly intense lines between 2933 and $2981 \mathrm{~cm}^{-1}$. Although, the experimental spectrum has an intense, broad feature centered around $2930 \mathrm{~cm}^{-1}$, the structure in this region was not attributed to the diradical, ${ }^{\text {6a }}$ perhaps because this part of the $\mathrm{C}-\mathrm{H}$ stretching region was not materially changed between the starting azo compound and the product. Peaks near 2820, 2880, and $3120 \mathrm{~cm}^{-1}$ were assigned to II; we do not predict any fundamentals within $100 \mathrm{~cm}^{-1}$ of these values, suggesting that these lines may be due to a species other than II. However, we
(10) Du, P.; Hrovat, D. A.; Borden, W. T. J. Am. Chem. Soc. 1986, 108, 8086.

Table I. Theoretical Frequencies and Normalized IR Intensities of II and III ${ }^{a}$

| II |  |  | III |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| symmetry | frequency ( $\mathrm{cm}^{-1}$ ) | intensity | symmetry | frequency ( $\mathrm{cm}^{-1}$ ) | intensity |
| $A$ | 3017.0 | 0.024 | $A_{1}$ | 3035.6 | 0.043 |
|  | 2975.9 | 0.166 |  | 2999.0 | 0.030 |
|  | 2962.5 | 0.091 |  | 2958.3 | 0.115 |
|  | 2940.7 | 0.229 |  | 2954.1 | 0.180 |
|  | 2937.9 | 0.071 |  | 1496.7 | 0.043 |
|  | 1493.0 | 0.037 |  | 1444.1 | 0.017 |
|  | 1476.8 | 0.028 |  | 1360.6 | 0 |
|  | 1386.8 | 0.006 |  | 1327.2 | 0.001 |
|  | 1376.7 | 0.002 |  | 1165.3 | 0.027 |
|  | 1297.8 | 0.022 |  | 987.4 | 0.007 |
|  | 1204.7 | 0.027 |  | 870.0 | 0.028 |
|  | 1142.6 | 0.012 |  | 751.9 | 0 |
|  | 1073.6 | 0.002 |  | 671.4 | 0 |
|  | 991.7 | 0.002 |  | 508.1 | 0 |
|  | 960.2 | 0.002 |  | 296.6 | 0.001 |
|  | 815.0 | 0.006 | $A_{2}$ | 1126.4 | 0 |
|  | 772.1 | 0.088 |  | 766.8 | 0 |
|  | 693.5 | 0 |  | 673.6 | 0 |
|  | 645.4 | 0.001 |  | 637.6 | 0 |
|  | 630.5 | 0.003 |  | 141.4 | 0 |
|  | 541.1 | 0.003 | $B_{1}$ | 2980.5 | 0.133 |
|  | 515.2 | 0.003 |  | 873.5 | 0.057 |
|  | 328.0 | 0 |  | 768.7 | 1.000 |
|  | 228.0 | 0.059 |  | 676.8 | 0.219 |
|  | 74.6 | 0 |  | 555.9 | 0.016 |
| $B$ | 3011.5 | 0.180 |  | 345.2 | 0.067 |
|  | 2981.5 | 0.339 |  | 172.6 | 0.063 |
|  | 2964.4 | 0.266 | $B_{2}$ | 3030.4 | 0.125 |
|  | 2937.4 | 0.175 |  | 2994.5 | 0.149 |
|  | 2933.0 | 0.058 |  | 2953.2 | 0.064 |
|  | 1479.5 | 0.030 |  | 1476.3 | 0.015 |
|  | 1471.5 | 0.028 |  | 1345.7 | 0.026 |
|  | 1388.6 | 0.009 |  | 1276.3 | 0.029 |
|  | 1356.7 | 0.015 |  | 1214.0 | 0.028 |
|  | 1246.0 | 0.009 |  | 1152.5 | 0.001 |
|  | 1200.7 | 0.001 |  | 1001.5 | 0.016 |
|  | 1163.1 | 0.003 |  | 886.5 | 0.021 |
|  | 1039.8 | 0.024 |  | 579.1 | 0.001 |
|  | 899.4 | 0.014 |  | 409.4 | 0.001 |
|  | 841.5 | 0.020 |  |  |  |
|  | 772.0 | 1.000 |  |  |  |
|  | 716.7 | 0.052 |  |  |  |
|  | 605.2 | 0.079 |  |  |  |
|  | 550.1 | 0.017 |  |  |  |
|  | 457.0 | 0.019 |  |  |  |
|  | 399.5 | 0.002 |  |  |  |
|  | 312.8 | 0.083 |  |  |  |
|  | 161.8 | 0.014 |  |  |  |

${ }^{a}$ II and III denote 2,3-dimethylenecyclohexa-1,3-diene and 2,3-di-methylenecyclopenta-1,3-diene, respectively. The intensities are normalized to the most intense peaks at $772.0 \mathrm{~cm}^{-1}$ (II) and $768.7 \mathrm{~cm}^{-1}$ (III).
cannot rule out the possibility that these lines might be due to overtones or that the UHF procedure, even with the "scaling" of the frequencies could have errors this large when applied to a diradical.


Figure 2. The calculated and experimental infrared spectra of 2,3 -di-methylenecyclohexa-1,3-diene (II). The experimental spectrum is taken from ref 6a.

In conclusion, 3-21G UHF calculations predict a $25.0^{\circ}$ deviation of the triplet diradical II from planarity. This value lies between those of tetramethyleneethane I and its five-membered ring a nalogue III, which are calculated to deviate by $44.9^{\circ 4}$ and $0^{\circ}$ from planarity. Ab initio calculations of the normal mode vibrational frequencies were also performed, leading us to question the assignments of the structure between 2820 and $3110 \mathrm{~cm}^{-1}$ in the infrared spectrum. ${ }^{6 a}$

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Supplementary Material Available: Complete tables of the calculated geometries of II and III are available upon request from the authors.


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